

Neutron/gamma-ray discrimination through measures of fit

Amiri, M., Prenosil, V. & Cvachovec, F.

Author post-print (accepted) deposited by Coventry University's Repository

Original citation & hyperlink:

Amiri, M, Prenosil, V & Cvachovec, F 2015, Neutron/gamma-ray discrimination through measures of fit. in 4th International Conference on Advancements in Nuclear Instrumentation Measurement Methods and their Applications (ANIMMA). IEEE, 4th International Conference on Advancements in Nuclear Instrumentation Measurement Methods and their Applications (ANIMMA), Lisbon, Portugal, 20/04/15.
<https://dx.doi.org/10.1109/ANIMMA.2015.7465553>

DOI 10.1109/ANIMMA.2015.7465553

ISBN 978-1-4799-9918-7

Publisher: IEEE

© 2015 IEEE. Personal use of this material is permitted. Permission from IEEE must be obtained for all other uses, in any current or future media, including reprinting/republishing this material for advertising or promotional purposes, creating new collective works, for resale or redistribution to servers or lists, or reuse of any copyrighted component of this work in other works.

Copyright © and Moral Rights are retained by the author(s) and/ or other copyright owners. A copy can be downloaded for personal non-commercial research or study, without prior permission or charge. This item cannot be reproduced or quoted extensively from without first obtaining permission in writing from the copyright holder(s). The content must not be changed in any way or sold commercially in any format or medium without the formal permission of the copyright holders.

This document is the author's post-print version, incorporating any revisions agreed during the peer-review process. Some differences between the published version and this version may remain and you are advised to consult the published version if you wish to cite from it.

Neutron/Gamma-Ray Discrimination through Measures of Fit

Moslem Amiri, Václav Přenosil, and František Cvachovec

Abstract—Statistical tests and their underlying measures of fit can be utilized to separate neutron/gamma-ray pulses in a mixed radiation field. In this article, first the application of a sample statistical test is explained. Fit measurement-based methods require true pulse shapes to be used as reference for discrimination. This requirement makes practical implementation of these methods difficult; typically another discrimination approach should be employed to capture samples of neutrons and gamma-rays before running the fit-based technique. In this article, we also propose a technique to eliminate this requirement. These approaches are applied to several sets of mixed neutron and gamma-ray pulses obtained through different digitizers using stilbene scintillator in order to analyze them and measure their discrimination quality.

Index Terms—Counter/discriminator, statistical test, measures of fit, neutron spectroscopy, organic scintillator.

I. INTRODUCTION

ORGANIC scintillation detectors are widely used for neutron detection and spectroscopy. The primary reaction producing neutron field and scattering reactions of neutrons with materials in the environment lead to the production of γ -ray as a background radiation. The main problem in neutron detection is the discrimination of neutrons from the γ -rays. A composite curve often comprises a fast and a slow component of scintillation. The long-lived slow component often reveals the nature of the particle striking the detector. This fact is usually used to separate different kinds of particles; this process is called pulse shape discrimination (PSD) [1]. PSD techniques are mainly used to discriminate neutrons from the γ -rays.

Stilbene and NE-213 organic scintillators are suitable for neutron spectroscopy; despite their low light output per unit energy, this light output induced by charged protons can be easily distinguished from electrons/photons. Therefore, these scintillators produce good results using PSD methods.

Popular analog PSD techniques, including rise-time inspection, zero-crossing method, and charge comparison, were classically used for n/γ -ray discrimination [2]. These analog techniques make acceptable n/γ -ray discrimination. However, the availability of precise and fast digitizers and various PSD algorithms have made it possible to discriminate the radiations

better using digital techniques. Pulse rise-time algorithm and charge comparison are two successful digital PSD methods.

Measures of fit of statistical tests can be easily used to discriminate neutrons and γ -rays. The underlying operations in these tests are generally the same. In this article, the application of a statistical test, namely Pearson's chi-squared test, is utilized to separate n/γ -ray pulses in a mixed radiation field. Since fit measurement-based methods require two reference pulse shapes to be preliminarily available for discrimination, an update to the original implementation is introduced later in this article to eliminate this requirement. The sampled data of mixed neutron and γ -ray pulses are obtained through two differently-featured digitizers (explained in Section II) which differ mainly in their sampling rate and output quantization level resolution. This helps find the effect of resolution and sampling frequency of the digitizers on the quality of the discrimination result for the methods discussed in this article. Every experiment is carried out under the same experimental conditions, using 100,000 pulses of mixed neutron and photon signals. For this work, the field consists of mostly γ -rays and some neutrons.

Various techniques, applied to data obtained from different digitizer types and settings, are compared by using the Figure of Merit (FoM) for the n/γ -ray discrimination, defined as:

$$FoM = \frac{S}{FWHM_n + FWHM_\gamma} \quad (1)$$

where S is the separation between the peaks of the two events, $FWHM_\gamma$ is the full-width half-maximum (FWHM) of the spread of events classified as γ -rays and $FWHM_n$ is the FWHM of the spread in the neutron peak [3]. FWHMs are calculated using the Gaussian fits to the neutron and γ -ray events on experimental distribution plot.

II. EXPERIMENTAL SETUP

For this work, stilbene scintillation detector was used with 45x45 crystal, and the neutron-gamma radiation source used was $^{252}\text{Cf(sf)}$. A typical scintillation detector consists of a scintillator and a photomultiplier. The latter is employed to change weak light signals impinging to photocathode (generated by the scintillator) into electric impulses. We used the photomultiplier RCA7265 [4] throughout these experiments. The block diagram of our digital apparatus is shown in Fig. 1.

A preamplifier is selected so as to match the detector output impedance. Two variants of the anode load resistance were tested in conjunction with the organic scintillation detectors. In the first variant, a load resistance of $40\text{ k}\Omega$ was used. A

Manuscript received May 30, 2015. This work was supported by Technology Agency of the Czech Republic under contract No. TA01011383/2011.

M. Amiri is with the Faculty of Informatics, Masaryk University, Botanická 68a, 602 00 Brno, Czech Republic (e-mail: amiri@mail.muni.cz).

V. Přenosil is with the Faculty of Informatics, Masaryk University, Botanická 68a, 602 00 Brno, Czech Republic (e-mail: prenosil@fi.muni.cz).

F. Cvachovec is with the Faculty of Military Technology, University of Defence, Kounicova 156/65, 662 10 Brno, Czech Republic (e-mail: frantisek.cvachovec@unob.cz).

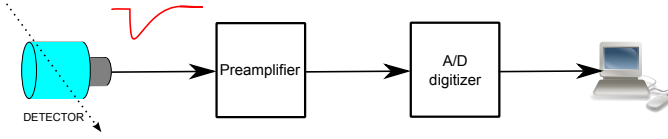


Fig. 1. Block diagram of a digital two-parameter analyzer.

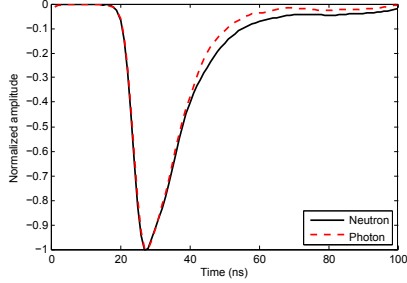


Fig. 2. Comparison of a sample smoothed neutron with a sample smoothed photon. These signals are obtained from the stilbene scintillator.

preamplifier matched it to the coaxial cable whose characteristic impedance was 50Ω . In this case, the different waveforms of the neutron and photon pulses can be detected in the voltage pulse leading edge. If the magnitude of the load resistance is selected to be close to the characteristic impedance of the coaxial cable, which is 50Ω , the different shapes of the neutron/photon pulses will appear to take effect during the decay time. In this case, no preamplifier is necessary. The latter option was employed here.

Two commercially available Agilent digitizers were used to digitize the output pulses: Acqiris DP210 with 8-bit resolution and set at 1 and 2 GS/s, and Acqiris DC440 with 12-bit resolution and set at 250 and 420 MS/s. While real-time digitizers are also employed in industry today, we used these specific commercial digitizers to study the effects of their various data resolution and sampling frequency features on digital processing.

III. NEUTRON AND PHOTON SIGNALS

A sample smoothed neutron is compared with a sample smoothed photon pulse in Fig. 2. These signals are obtained from the stilbene scintillator. As seen in this Figure, these signals are composed of a leading and a trailing edge. The leading edges could not be exploited for discrimination purposes. On the other hand, the trailing edge of the neutron signal has higher rise time than that of the photon signal. This property could be used to separate these two radiations. It is sufficient to apply a discrimination technique only to the differing segments of the two signal types.

IV. MEASURES OF FIT

Measures of fit of statistical tests can be employed to discriminate neutrons and γ -rays. The underlying operations in these tests are generally the same. In [5], [6], a similar approach, called "model pulse algorithm," is adopted. In the following, we propose a simpler and more practical approach to separating radiations.

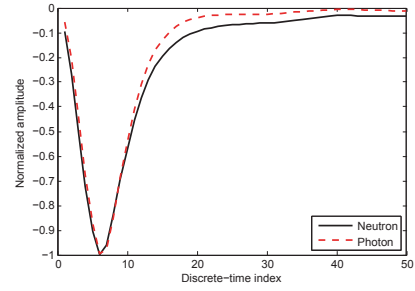


Fig. 3. Segments of neutron and γ -ray true pulse shapes, obtained by averaging multiple normalized neutron and γ -ray pulses. The original pulses are obtained from DC440 digitizer (12-bit resolution, 420 MS/s). The same segment of every unknown pulse from the data set will be used for the measurements.

First, we construct two normalized reference pulse shapes, called true pulse shapes, one for neutrons and the other for γ -rays. The γ -ray true pulse shape can be obtained by averaging some normalized signals captured from different pure γ -ray sources, and the neutron true pulse shape is obtained by averaging some normalized neutron pulses captured by discrimination through some method, e.g. integration method (the latter approach can be used to build the γ -ray true pulse shape too). For this work, the starting point of the two normalized true pulse shapes is where the amplitude of the leading edge is 10% fraction of peak-amplitude, and the end point is some specific time after the starting point such that this interval includes the differing tail segment between neutrons and γ -rays. The same approach used to cut a segment from the true pulse shapes will also be used to cut a segment from each normalized unknown pulse from a mixed n/ γ -ray field. Although the starting point could be set to be the peak of a pulse since the leading edges of neutron and γ -ray pulses almost map on each other, the inclusion of the leading edge (from 10% level to the peak) will make the calculations more precise. This segment is shown in Fig. 3, and used throughout our experiments.

Second, a statistical test is applied to separate n/ γ -ray pulses in a mixed radiation field. In this article, Pearson's chi-squared test is utilized as an example statistical test. By applying the test, two values are obtained by measuring the fit of the captured pulse and the two true pulse shapes (χ_n^2 , χ_γ^2). The measurements are done only within a specific segment of the pulses and the true pulse shapes, as explained above. The lower resulting value assigns the particle type. This approach can be extended to the other statistical tests available. The distribution of the (χ_n^2 , χ_γ^2) values provides an excellent discrimination of neutrons and γ -rays.

The equation to implement Pearson's chi-squared test for n/ γ -ray discrimination is:

$$\chi^2 = \sum_{i=1}^n \frac{(O_i - E_i)^2}{E_i} \quad (2)$$

where O_i is the observed value and E_i is the expected value. For discrimination application, O_i and E_i are the corresponding points in the sequences of captured pulse and the true pulse shapes respectively. Two χ^2 values are calculated; one

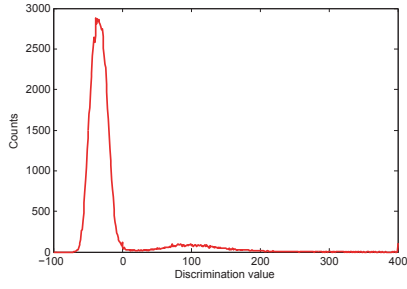


Fig. 4. Discrimination of γ -ray and neutron signals, applying the statistical test approach. The pulses are obtained using DP210 digitizer (8-bit resolution, 1 GS/s).

TABLE I

FOMS AND COUNTS OF THE PULSES OBTAINED FROM DC440 AND DP210 DIGITIZERS UNDER DIFFERENT SAMPLING RATES. STATISTICAL TEST APPROACH IS APPLIED TO THE DATA.

Data format	FoM	Neutron counts	Photon counts
12-bit, 250 MS/s	1.10	9876	90124
12-bit, 420 MS/s	1.02	9795	90205
8-bit, 1 GS/s	1.02	10076	89924
8-bit, 2 GS/s	1.09	9920	90080

between the input pulse and the neutron true pulse (χ_n^2), and the other between the input pulse and the γ -ray true pulse (χ_γ^2). The smaller value between these two will determine the input pulse type. Hence this method could be used to count the number of pulses from each category. The subtraction of the two resulting χ^2 values for each input pulse could be used to calculate the FoM for that set of pulse data. This subtraction will scatter neutrons and γ -rays on two different sides of the zero base line of the plot. Fig. 4 illustrates the discrimination plot for the case of DP210 digitizer (8-bit resolution) with 1 GS/s sampling rate. As seen in this Figure, neutrons have positive and γ -rays have negative discrimination values in this experiment.

FoMs and n/ γ -ray counts for various data sets with different resolutions and frequency rates are shown in Table I. The results are solid; this method discriminates well irrespective of the digitizer features.

V. IMPROVED ALGORITHM

In the preceding technique, multiple detected neutrons and γ -rays are needed to build the fixed true pulse shapes to be used throughout the experiment. This requires the application of another discrimination technique prior to the statistical test in order to build the true pulse shapes. It is possible to build the true pulse shapes without the preliminary separation technique application, as will be explained in this section.

A simple method to build the two reference pulses is to pick two random pulses (e.g. the first two pulses at the beginning of the experiment) regarding each as representing a true pulse shape. The identity of these pulses are unknown and they may be even of the same type. The statistical test technique is applied to the rest of the pulses, using the two selected reference pulses. Every input pulse is categorized under one of the reference pulses which produces the lower resulting value, and the reference pulse is updated with that input pulse. Using

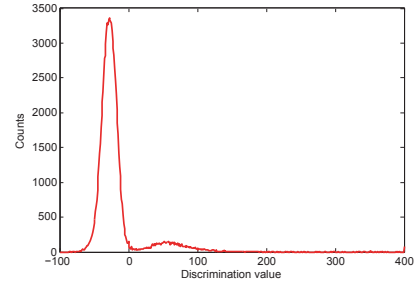


Fig. 5. Discrimination of γ -ray and neutron signals, applying the improved statistical test approach. The pulses are obtained using DP210 digitizer (8-bit resolution, 1 GS/s).

TABLE II

FOMS AND COUNTS OF THE PULSES OBTAINED FROM DC440 AND DP210 DIGITIZERS UNDER DIFFERENT SAMPLING RATES. IMPROVED STATISTICAL TEST APPROACH IS APPLIED TO THE DATA.

Data format	FoM	Neutron counts	Photon counts
12-bit, 250 MS/s	0.78	22780	77220
12-bit, 420 MS/s	1.03	9496	90504
8-bit, 1 GS/s	0.98	9817	90183
8-bit, 2 GS/s	0.84	10289	89711

this strategy, one reference sequence will be gradually moved toward the neutron and the other toward a γ -ray pulse shape.

Instead of the initial random selection of the two reference pulses, a better method is to average multiple random pulses and build up two new sequences by roughly shaping the resulting averaged sequence, once as neutron and once as γ -ray (it is not necessary to be precise). Then the statistical test approach is applied to every pulse, using the two rough references. The references are updated as the experiment goes on.

Using either of the methods explained above in this section, the pulses will be divided and grouped with either of the two reference sequences. If not detected, at the end of the experiment, when the reference pulses are completely refined and shaped, a simple test will reveal the identity of the two reference sequences, hence the identity of the pulses grouped with each reference pulse. For example, since the reference and input pulses are normalized throughout the experiment, if the samples of each reference pulse are summed up, the one with the lower sum will be γ -ray and the other one will be neutron. Fig. 5 illustrates the experimental distribution plot of neutrons and γ -rays for the data obtained from DP210 digitizer with 8-bit resolution and set at 1 GS/s frequency rate, when this method is applied.

FoMs and n/ γ -ray counts for various data sets with different resolutions and frequency rates are shown in Table II. The same data sets used for the experiments of Section IV are used here for comparison purposes. Although this method does not require the true pulse shapes prior to the experiment, it does not produce acceptable FoMs or accurate pulse counts. The lowest discrimination quality and the most inaccurate pulse count is when the sampling rate of the digitizer is low.

TABLE III
FoMs of PGA METHOD FOR THE PULSES OBTAINED FROM VARIOUS DIGITIZERS.

Digitizer	8-bit, 1 GS	8-bit, 2 GS	12-bit, 250 MS	12-bit, 420 MS
FoM	0.88	0.91	0.94	1.00

VI. DISCUSSION

One factor affecting the FoM of a discrimination method is sampling rate of the digitizer. Neutron and photon signals contain frequency components up to 100 MHz [7]. Thus, based on Nyquist criterion, the minimum sampling frequency for neutron and photon signals should be greater than 200 MS/s. The estimation of the sampling rate impact on the separation quality of a method can be involved. For the approaches introduced in this article, as Tables I and II show, increasing from the low sampling rate of 250 MHz to 420 MHz or increasing from the high sampling rate of 1 GHz to 2 GHz does not necessarily improve the FoM.

The factor with a greater impact on discrimination quality is digitizer resolution. The quality of the digitizer output could be measured by signal-to-quantization noise ratio (SQNR). Since quantization errors of neutron and photon signals are almost uniformly distributed over the quantization interval, the following well-known equation [8] reliably estimates the quality of a b -bit digitizer output:

$$SQNR(dB) = 1.76 + 6.02b \quad (3)$$

Equation (3) implies that SQNR increases approximately 6 dB for every bit added to the digitizer word length. This relationship gives the number of bits required by an application to assure a given signal-to-noise ratio.

In order to verify the performance of the methods introduced in this article, we apply PGA method to the same pulse data sets as used for the methods in this paper. PGA method, introduced in [3], is recognized as an efficient n/γ -ray discrimination method; the gradient between the peak amplitude and the amplitude a specified time after the peak amplitude (called the discrimination amplitude) on the trailing edge of the pulses are compared and used as the discrimination factor. Fig. 6 illustrates the peak and discrimination amplitudes on neutron and photon signals. The gradient is calculated using

$$m = \frac{\Delta y}{\Delta t} = \frac{(y_p - y_d)}{(t_p - t_d)} \quad (4)$$

where m , y_p , y_d , t_p , and t_d are the gradient, the peak amplitude (which is a constant for normalized pulses), the discrimination amplitude, the time of peak amplitude occurrence, and the time of discrimination amplitude occurrence, respectively. For this work, we used some training pulses to locate the best discrimination amplitude, which occurred about 36 ns after the peak of the pulse. The FoMs obtained are listed in Table III. A comparison shows that the novel methods introduced here are either better or at least have the same discrimination quality as the PGA method does. Fig. 7 shows the best discrimination plot obtained by PGA method.

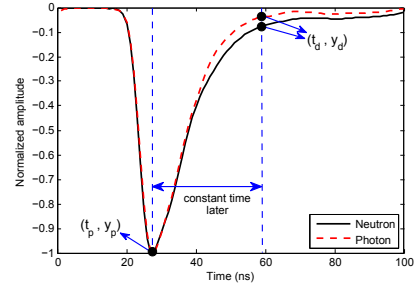


Fig. 6. The points on smoothed neutron and photon signals used in PGA discrimination method.

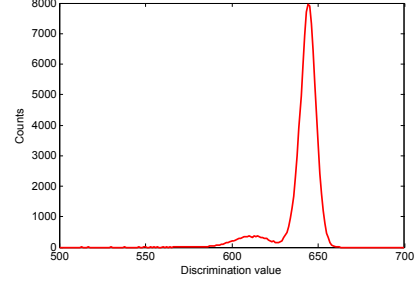


Fig. 7. Discrimination of photon and neutron signals, applying PGA method. The pulses are obtained using DC440 digitizer (12-bit resolution, 420 MS/s).

VII. CONCLUSION

In this article, we first introduced n/γ -ray discrimination through measures of fit, particularly using Pearson's chi-squared test. This approach counts and discriminates the pulses in a mixed environment very efficiently, however, it requires sample neutron and γ -ray pulses in advance of the experiment run to build the true pulse shapes. This requirement leads to the employment of another technique prior to the statistical test method to provide proper reference sequences. Since this may not be practically acceptable, an improvement to the original statistical test method is made in this article to build the true pulse shapes during run time. The true pulse shapes are gradually updated and used as the reference throughout the experiment. Two digitizers, each featuring a different resolution and each set at two different sampling rates, were used to observe the reaction of the methods to the data sampling conditions. The original approach (requiring true pulse shapes prior to the experiment) provides promising results while the improved approach is inefficient specially when applied to the data sampled at low rates.

REFERENCES

- [1] G. F. Knoll, *Radiation Detection and Measurement*, 3rd ed. New York: Wiley, 2000, ch. 8, p. 230.
- [2] G. Ranucci, "An analytical approach to the evaluation of the pulse shape discrimination properties of scintillators," *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment*, vol. 354, no. 2-3, pp. 389 – 399, 1995. [Online]. Available: <http://www.sciencedirect.com/science/article/pii/0168900294008868>
- [3] B. Mellow, M. Aspinall, R. Mackin, M. Joyce, and A. Peyton, "Digital discrimination of neutrons and g-rays in liquid scintillators using pulse gradient analysis," *Nuclear Instruments and Methods in Physics Research Section A*, vol. 578, no. 1, pp. 191–197, 2007.

- [4] "Rca 7265 photomultiplier tube 2" 14-stage s-20," <http://www.hofstragroup.com/product/rca-7265-photomultiplier-tube-2-14-stage-s-20/>, accessed: 2014-07-15.
- [5] S. Marrone, D. Cano-Ott, N. Colonna, C. Domingo, F. Gramegna, E. Gonzalez, F. Gunsing, M. Heil, F. Käppeler, P. Mastinu, P. Milazzo, T. Papaevangelou, P. Pavlopoulos, R. Plag, R. Reifarth, G. Tagliente, J. Tain, and K. Wisshak, "Pulse shape analysis of liquid scintillators for neutron studies," *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment*, vol. 490, no. 1-2, pp. 299–307, 2002. [Online]. Available: <http://www.sciencedirect.com/science/article/pii/S016890020201063X>
- [6] C. Guerrero, D. Cano-Ott, M. Fernández-Ordóñez, E. González-Romero, T. Martínez, and D. Villamarín, "Analysis of the BC501A neutron detector signals using the true pulse shape," *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment*, vol. 597, no. 2-3, pp. 212–218, 2008. [Online]. Available: <http://www.sciencedirect.com/science/article/pii/S0168900208014113>
- [7] F. Belli, B. Esposito, D. Marocco, and M. Riva, "A study on the pulse height resolution of organic scintillator digitized pulses," *Fusion Engineering and Design*, vol. 88, no. 68, pp. 1271 – 1275, 2013, proceedings of the 27th Symposium On Fusion Technology (SOFT-27); Liège, Belgium, September 24-28, 2012. [Online]. Available: <http://www.sciencedirect.com/science/article/pii/S092037961200587X>
- [8] J. G. Proakis and D. G. Manolakis, *Digital Signal Processing, Fourth Edition*. Upper Saddle River, New Jersey 07458: Prentice Hall, 2006.